

Anomalies and $L_\mu - L_\tau$

Julian Heeck

Anomalies 2021

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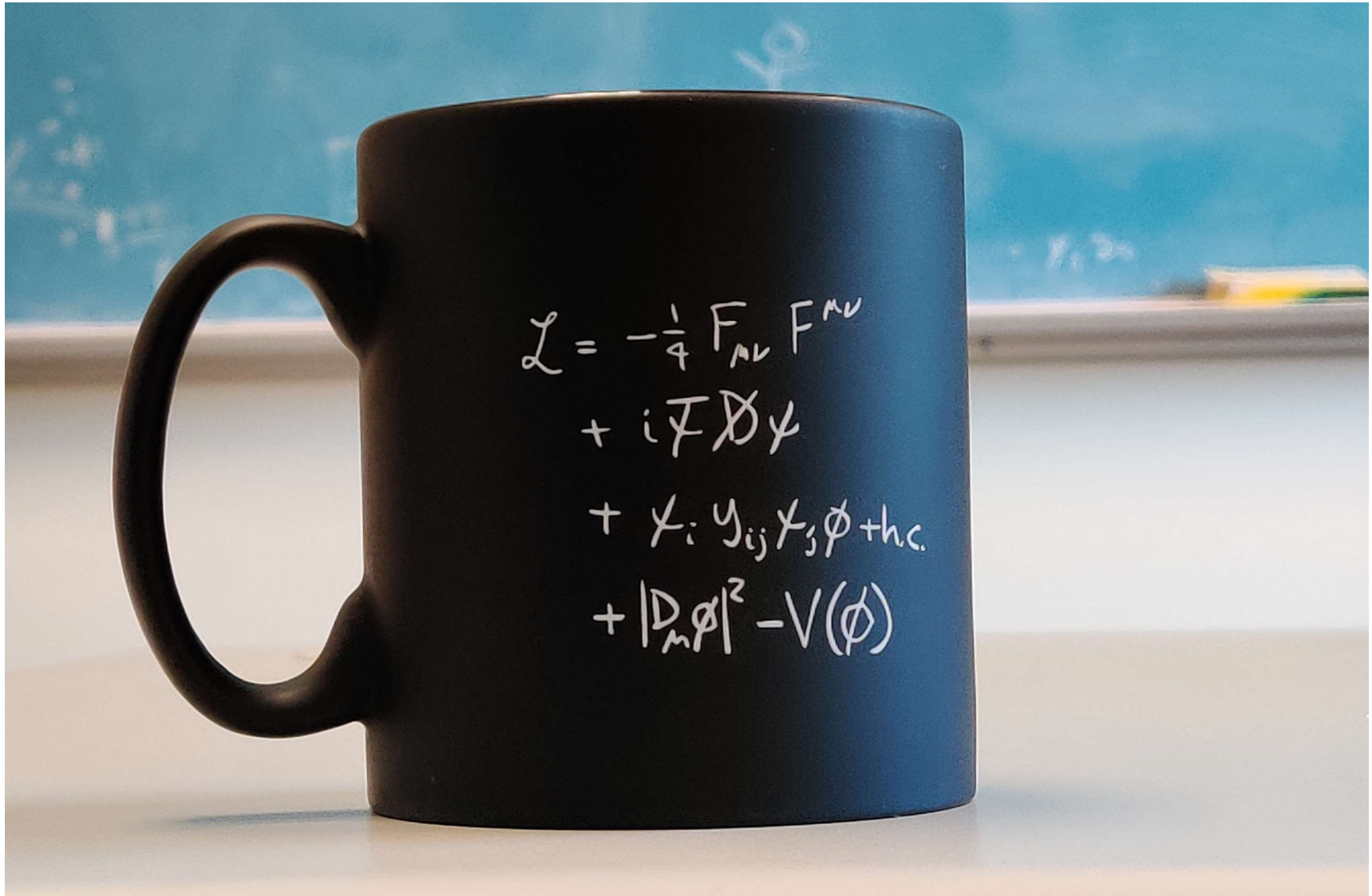
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Elementary particles

| | | | | | |
|--------------------|--|--|--|--|---|
| | <p>mass → $\approx 2.3 \text{ MeV}/c^2$</p> <p>charge → $2/3$</p> <p>spin → $1/2$</p> <p>u</p> <p>up</p> | <p>mass → $\approx 1.275 \text{ GeV}/c^2$</p> <p>charge → $2/3$</p> <p>spin → $1/2$</p> <p>c</p> <p>charm</p> | <p>mass → $\approx 173.07 \text{ GeV}/c^2$</p> <p>charge → $2/3$</p> <p>spin → $1/2$</p> <p>t</p> <p>top</p> | <p>mass → 0</p> <p>charge → 0</p> <p>spin → 1</p> <p>g</p> <p>gluon</p> | <p>mass → $\approx 126 \text{ GeV}/c^2$</p> <p>charge → 0</p> <p>spin → 0</p> <p>H</p> <p>Higgs boson</p> |
| BARYONS/ QUARKS | <p>mass → $\approx 4.8 \text{ MeV}/c^2$</p> <p>charge → $-1/3$</p> <p>spin → $1/2$</p> <p>d</p> <p>down</p> | <p>mass → $\approx 95 \text{ MeV}/c^2$</p> <p>charge → $-1/3$</p> <p>spin → $1/2$</p> <p>s</p> <p>strange</p> | <p>mass → $\approx 4.18 \text{ GeV}/c^2$</p> <p>charge → $-1/3$</p> <p>spin → $1/2$</p> <p>b</p> <p>bottom</p> | <p>mass → 0</p> <p>charge → 0</p> <p>spin → 1</p> <p>γ</p> <p>photon</p> | |
| | | <p>mass → $0.511 \text{ MeV}/c^2$</p> <p>charge → -1</p> <p>spin → $1/2$</p> <p>e</p> <p>electron</p> | <p>mass → $105.7 \text{ MeV}/c^2$</p> <p>charge → -1</p> <p>spin → $1/2$</p> <p>μ</p> <p>muon</p> | <p>mass → $1.777 \text{ GeV}/c^2$</p> <p>charge → -1</p> <p>spin → $1/2$</p> <p>τ</p> <p>tau</p> | <p>mass → $91.2 \text{ GeV}/c^2$</p> <p>charge → 0</p> <p>spin → 1</p> <p>Z</p> <p>Z boson</p> |
| LEPTONS | <p>mass → $< 2.2 \text{ eV}/c^2$</p> <p>charge → 0</p> <p>spin → $1/2$</p> <p>ν_e</p> <p>electron neutrino</p> | <p>mass → $< 0.17 \text{ MeV}/c^2$</p> <p>charge → 0</p> <p>spin → $1/2$</p> <p>ν_μ</p> <p>muon neutrino</p> | <p>mass → $< 15.5 \text{ MeV}/c^2$</p> <p>charge → 0</p> <p>spin → $1/2$</p> <p>ν_τ</p> <p>tau neutrino</p> | <p>mass → $80.4 \text{ GeV}/c^2$</p> <p>charge → ± 1</p> <p>spin → 1</p> <p>W</p> <p>W boson</p> | |

[wikipedia]

Interactions



Symmetries of the Standard Model

- Rephasing lepton and quark fields:

$$\begin{aligned} & U(1)_B \times U(1)_{L_e} \times U(1)_{L_\mu} \times U(1)_{L_\tau} \\ & = \\ & U(1)_{B+L} \times U(1)_{B-L} \times U(1)_{L_\mu - L_\tau} \times U(1)_{L_\mu + L_\tau - 2L_e} . \end{aligned}$$

Symmetries of the Standard Model

- Rephasing lepton and quark fields:

$$U(1)_B \times U(1)_{L_e} \times U(1)_{L_\mu} \times U(1)_{L_\tau}$$

=

$$\cancel{U(1)_{B+L}} \times U(1)_{B-L} \times U(1)_{L_\mu - L_\tau} \times U(1)_{L_\mu + L_\tau - 2L_e} \cdot$$

- **Broken** non-perturbatively, but unobservable. [['t Hooft, PRL '76](#)]
- True accidental global symmetry:

$$U(1)_{B-L} \times U(1)_{L_\mu - L_\tau} \times U(1)_{L_\mu + L_\tau - 2L_e} \cdot$$

Symmetries of the Standard Model

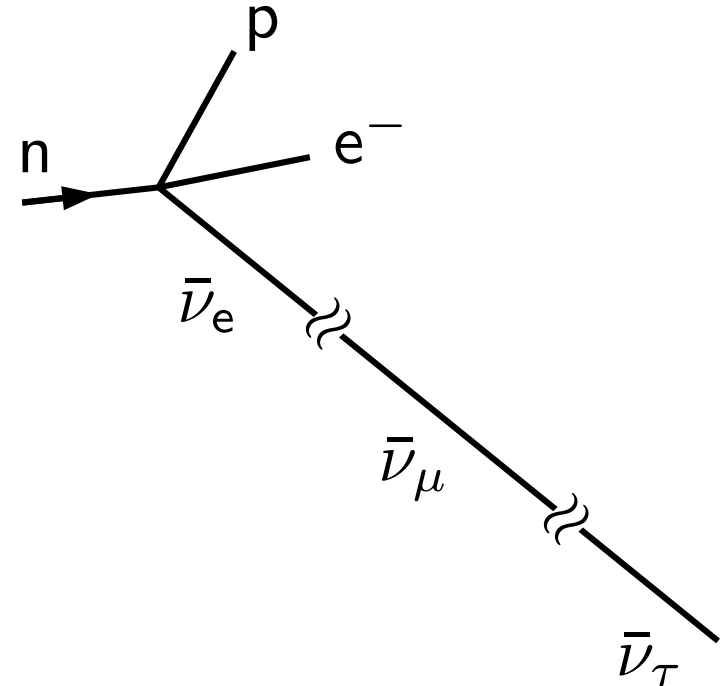
- True accidental global symmetry:

$$U(1)_{B-L} \times U(1)_{L_\mu - L_\tau} \times U(1)_{L_\mu + L_\tau - 2L_e} .$$

- A subgroup $U(1)_{L_\alpha - L_\beta}$ can be **gauged** without any new particles!
[Foot '91; He, Joshi, Lew, Volkas, '91]
- With 3 N_R , the entire $U(1)^3$ can be gauged. [Araki, JH, Kubo, '12]
 - Gauging global SM symmetry gives **neutrino masses**, but no mixing angles...

Neutrino oscillations = flavor violation

- Observations of $\nu_\alpha \rightarrow \nu_\beta$ prove that $M_\nu \neq 0$ and $U(1)_{L_\mu - L_\tau} \times U(1)_{L_\mu + L_\tau - 2L_e}$ is **broken!**
- Charged lepton flavor violation will occur, not clear how fast.
- B-L could still be **unbroken** if neutrinos are Dirac. [\[JH, 1408.6845\]](#)
 - B-L often broken to get seesaw.



Are the U(1) symmetries not interesting if they're broken?

Majorana mass matrix $\mathcal{M}_\nu = U \text{diag}(m_1, m_2, m_3) U^T$ in special cases:

- ① **Normal hierarchy** ($m_1 \simeq 0$) and best-fit values (phases zero):

$$\mathcal{M}_\nu \simeq \begin{pmatrix} 0.37 & 0.75 & 0.24 \\ \cdot & 2.47 & 2.11 \\ \cdot & \cdot & 2.99 \end{pmatrix} 10^{-2} \text{ eV} \sim \begin{pmatrix} 0 & 0 & 0 \\ 0 & \times & \times \\ 0 & \times & \times \end{pmatrix} \leftarrow L_e$$

- ② **Inverted hierarchy** ($m_3 \simeq 0$) and $\alpha = \pi/2$:

$$\mathcal{M}_\nu \simeq \begin{pmatrix} 1.84 & -3.11 & 3.22 \\ \cdot & -0.14 & 0.88 \\ \cdot & \cdot & -1.77 \end{pmatrix} 10^{-2} \text{ eV} \sim \begin{pmatrix} 0 & \times & \times \\ \times & 0 & 0 \\ \times & 0 & 0 \end{pmatrix} \leftarrow L_e - L_\mu - L_\tau$$

- ③ **Quasi-degenerate** ($m_{1,2,3} \simeq 1 \text{ eV}$) and $\beta = \pi/2$:

$$\mathcal{M}_\nu \simeq \begin{pmatrix} 0.96 & -0.20 & -0.22 \\ \cdot & 0.11 & -0.97 \\ \cdot & \cdot & -0.07 \end{pmatrix} \text{ eV} \sim \begin{pmatrix} \times & 0 & 0 \\ 0 & 0 & \times \\ 0 & \times & 0 \end{pmatrix} \leftarrow L_\mu - L_\tau$$

- Three interesting zeroth order approximations:

$$\mathcal{M}_\nu^{L_e} \sim \begin{pmatrix} 0 & 0 & 0 \\ 0 & \times & \times \\ 0 & \times & \times \end{pmatrix}, \quad \mathcal{M}_\nu^{L_e - L_\mu - L_\tau} \sim \begin{pmatrix} 0 & \times & \times \\ \times & 0 & 0 \\ \times & 0 & 0 \end{pmatrix}, \quad \mathcal{M}_\nu^{L_\mu - L_\tau} \sim \begin{pmatrix} \times & 0 & 0 \\ 0 & 0 & \times \\ 0 & \times & 0 \end{pmatrix}$$

[G. Branco, W. Grimus, L. Lavoura, *NPB* (1989); S. Choubey, W. Rodejohann, *EPJC* 40 (2005)]

- Normal, inverted, quasi-degenerate hierarchy might hint at “weakly broken” $B - 3L_e$, $B + 3(L_e - L_\mu - L_\tau)$, $L_\mu - L_\tau$



makes it anomaly free

[JH, Rodejohann, 1203.3117]

- Connection of $L_\mu - L_\tau$ to **neutrino mass anomaly**.
- Pressure on quasi-degenerate regime from cosmology bounds on Σm_ν . [Asai et al, 1811.07571]

Symmetries of the Standard Model

- True accidental global symmetry:

$$U(1)_{B-L} \times U(1)_{L_\mu - L_\tau} \times U(1)_{L_\mu + L_\tau - 2L_e} .$$

- With 3 N_R , the entire $U(1)^3$ can be gauged.
- First see the Z' with largest g' or smallest mass, coupled to

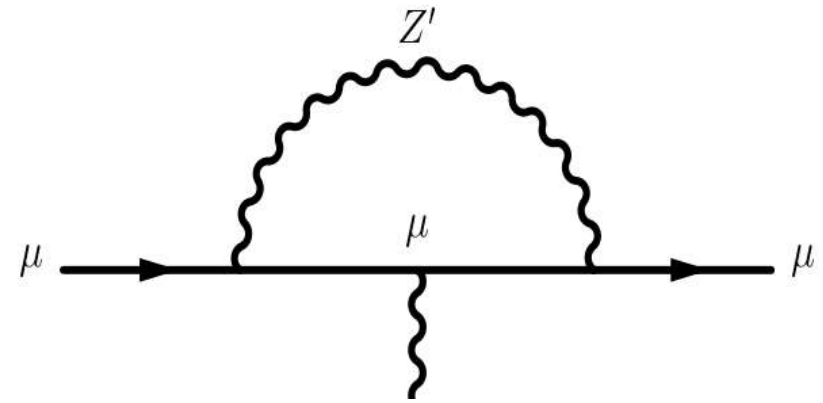
$$a(B - L) + b(L_\mu - L_\tau) + c(L_\mu + L_\tau - 2L_e) .$$

- Pheno dominated by electron coupling, except for $a=c=0!$

2001: $(g-2)_\mu$

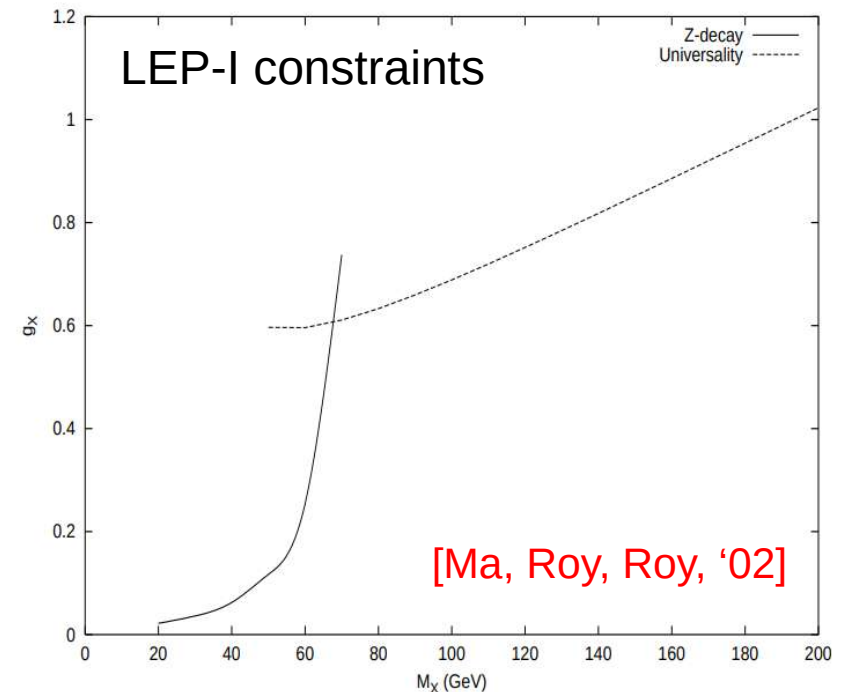
- Quickly realized that $L_\mu - L_\tau$ Z' can explain BNL result.

[Baek, Deshpande, He, Ko, '01; Ma, Roy, Roy, '02]



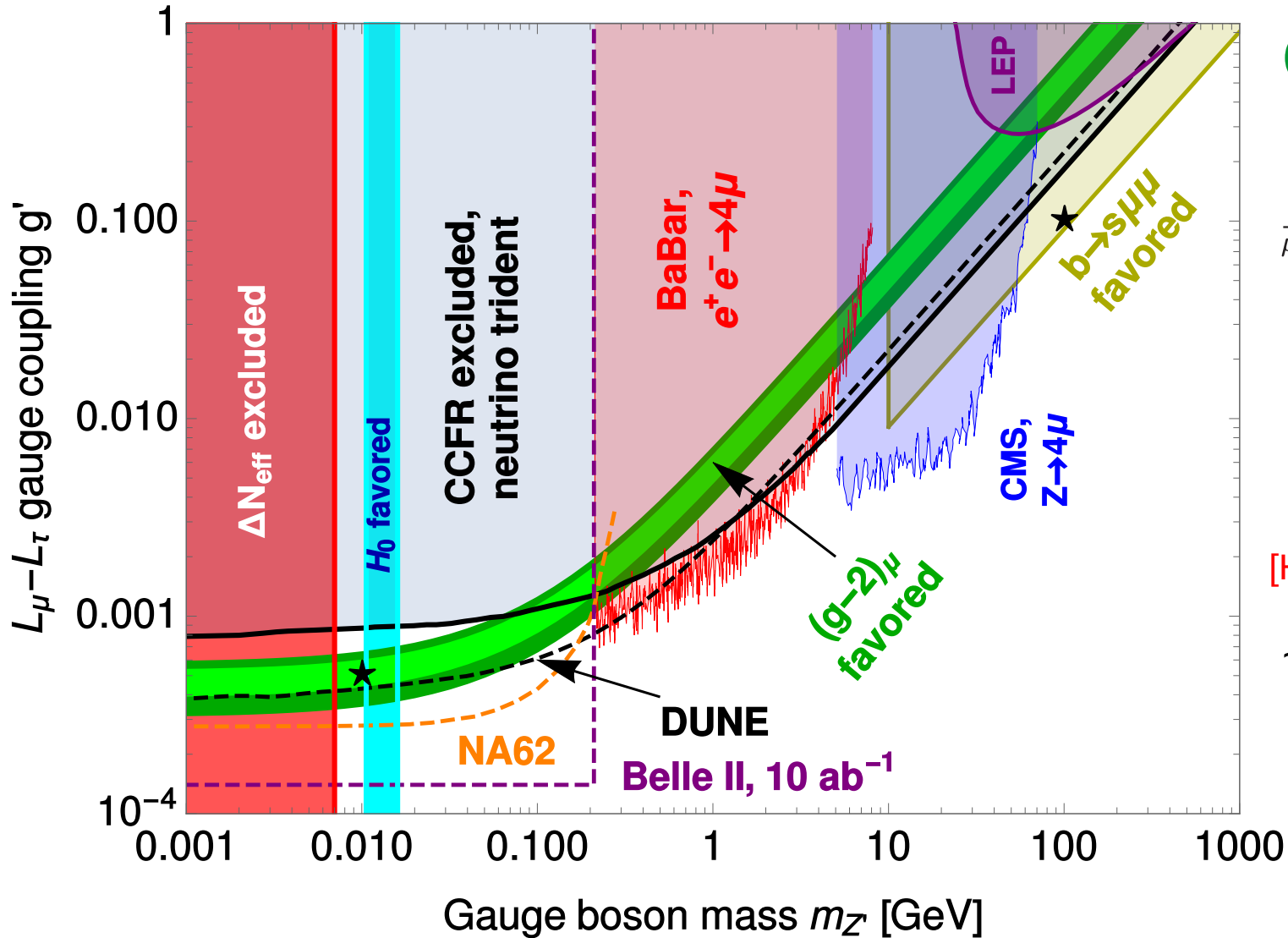
- Back then, *many* other simple models worked, but $L_\mu - L_\tau$ has survived for 20 years.

- While initially almost unconstrained due to vanishing first-generation couplings, it has now been pushed into a corner.

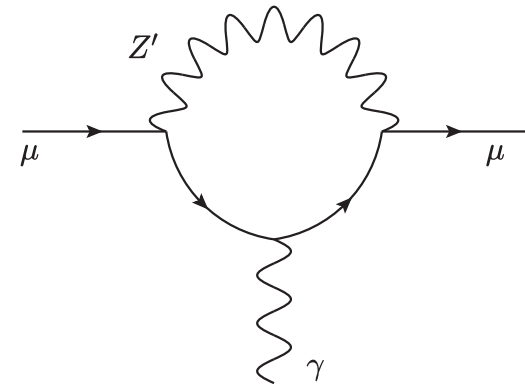


Current Z' constraints

[JH, Garani, 1906.10145]



$(g-2)_\mu$:

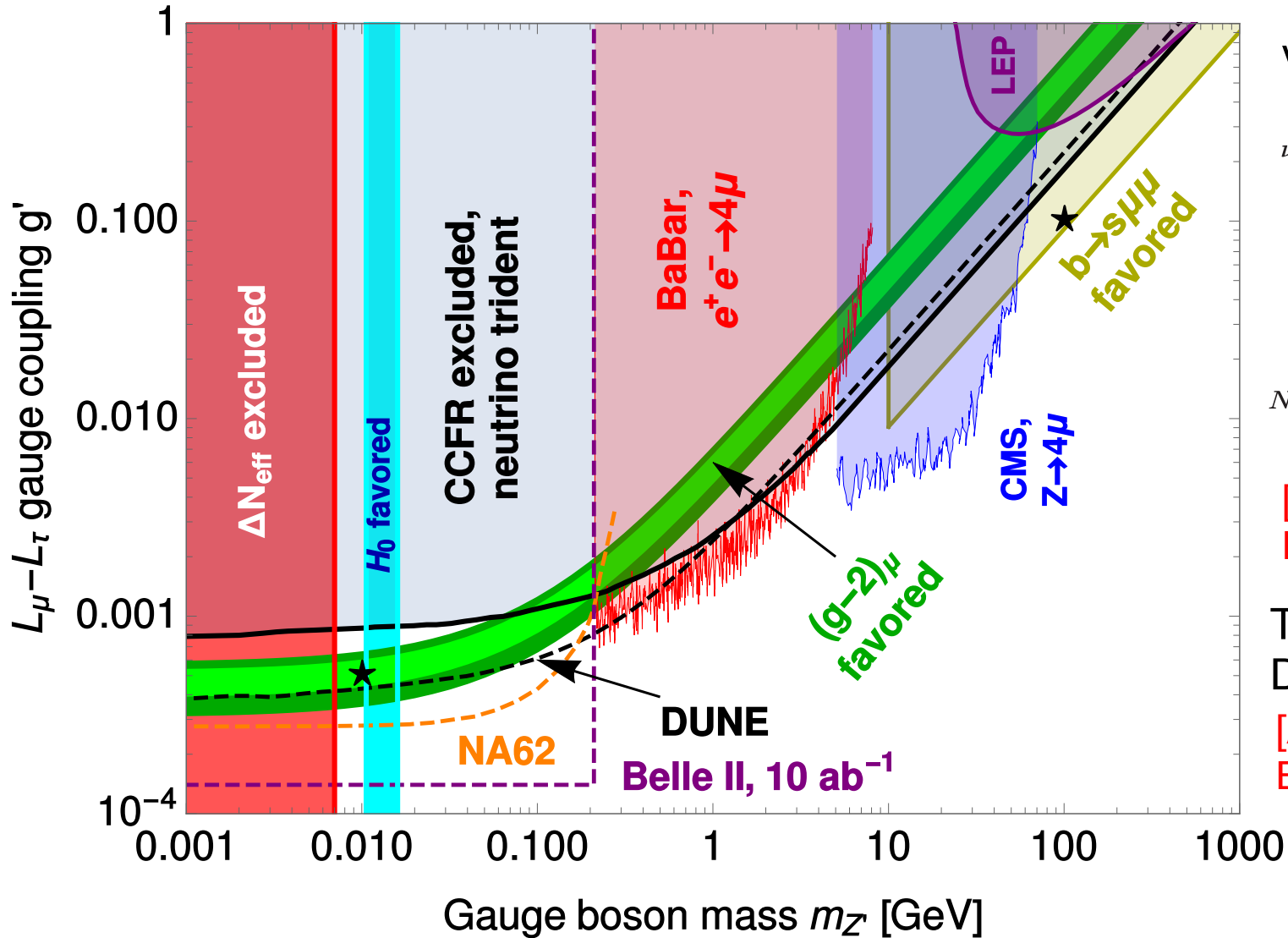


[He, Joshi, Lew, Volkas, '91]

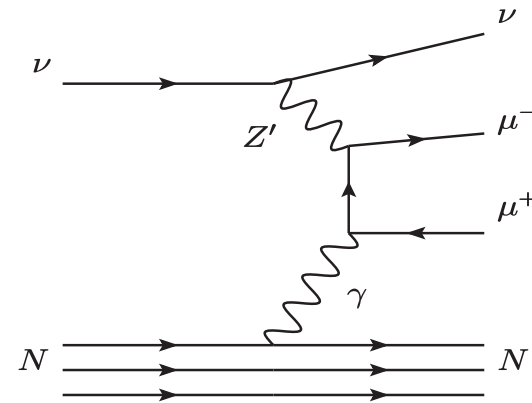
$\sim 4\sigma$ deviation!

Current Z' constraints

[JH, Garani, 1906.10145]



ν trident:



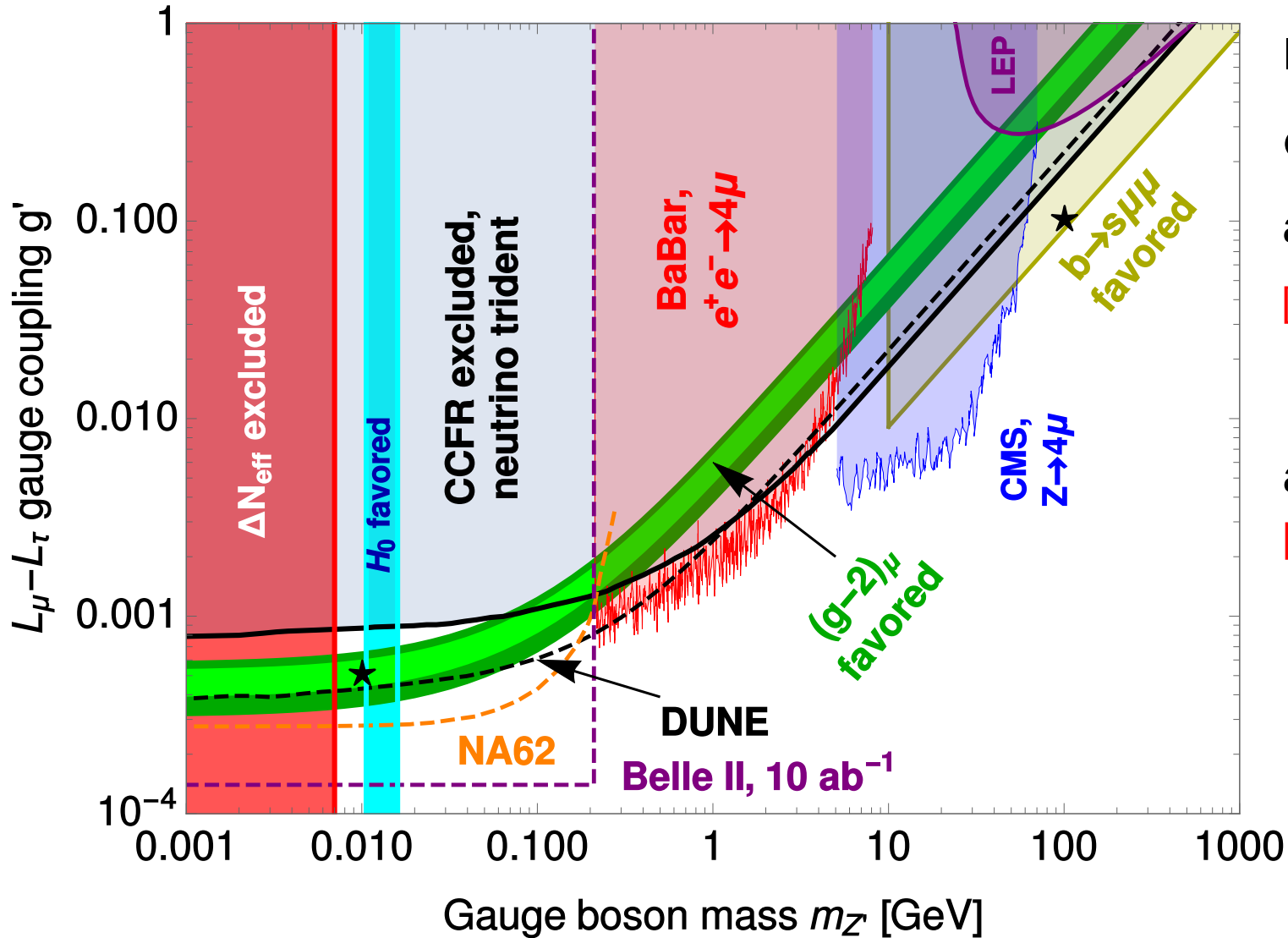
[Altmannshofer, Gori, Pospelov, Yavin, '14]

To be improved by DUNE!

[Altmannshofer++, '19; Ballett++, '19]

Current Z' constraints

[JH, Garani, 1906.10145]



Invisible Z' :

$$e^+e^- \rightarrow \mu^+\mu^-Z'$$

at Belle II.

[Jho++, 1904.1305]

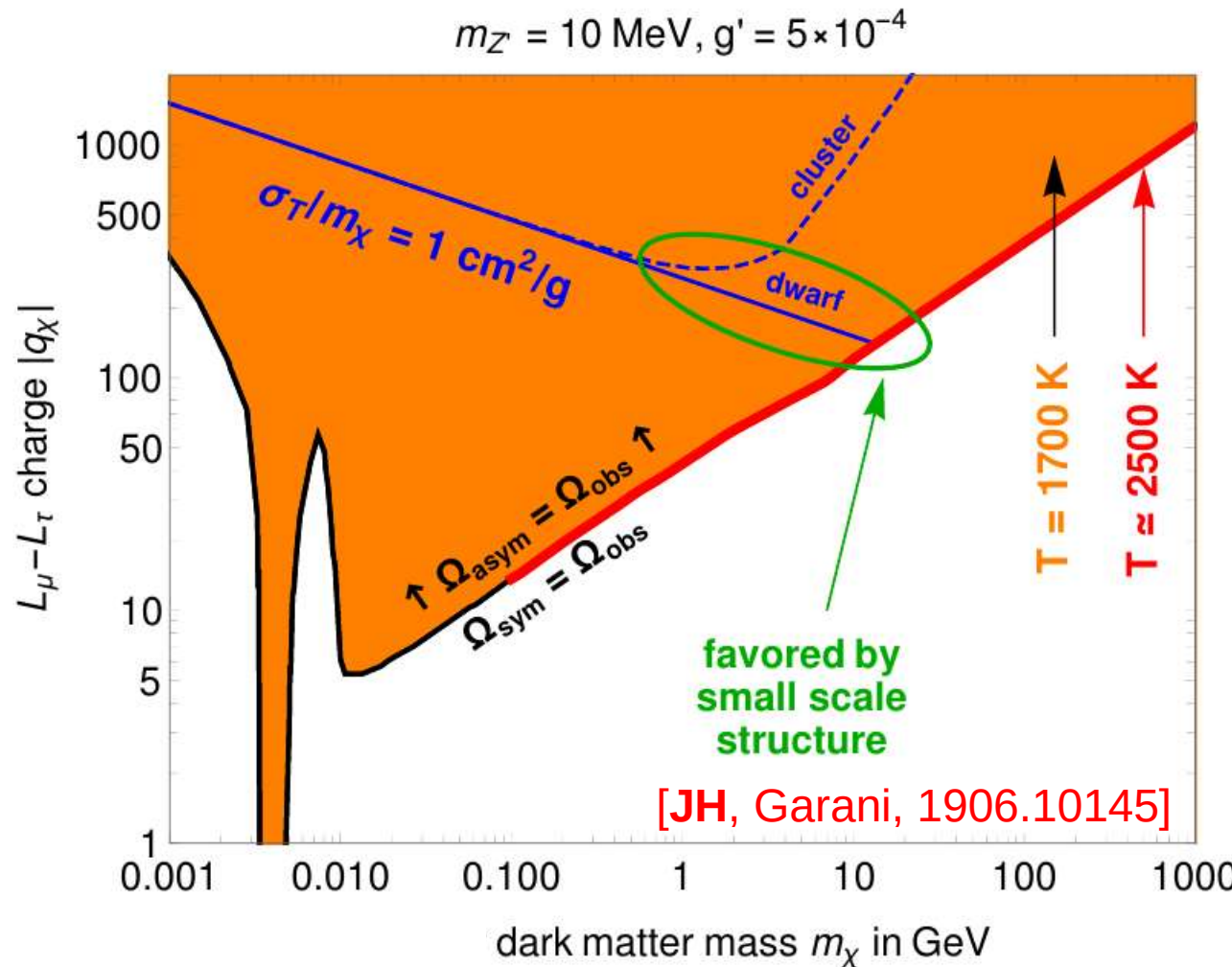
$$K \rightarrow \mu\nu Z'$$

at NA62.

[Krnjaic++, 1902.07715]

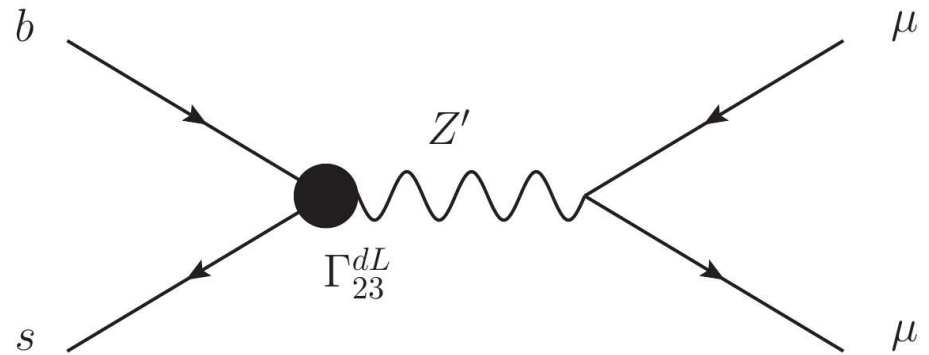
$L_\mu - L_\tau$ and g-2

- Pushed into corner $10 \text{ MeV} < m_{Z'} < 200 \text{ MeV}$, probed decisively in near future.
- Light Z' could be a mediator to a DM sector; can easily generate large **DM self-interaction** cross sections that solve small-scale structure problems.
- DM stability: $U(1)'$.



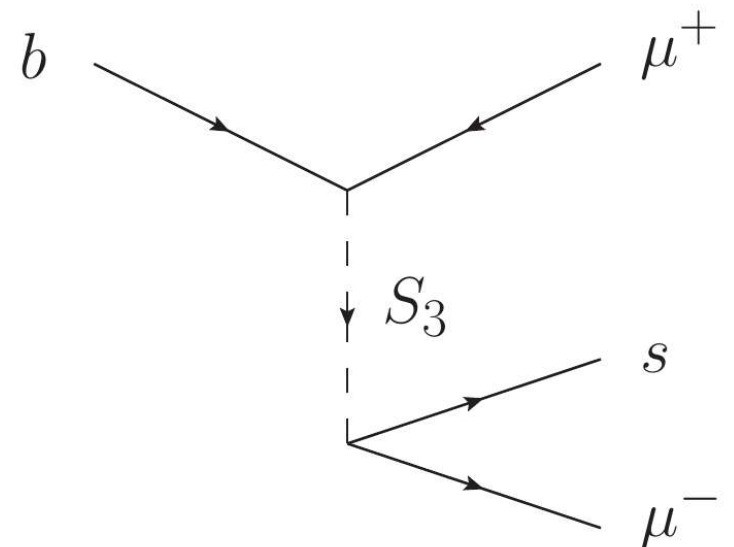
Non-minimal $L_\mu - L_\tau$

- Can induce additional Z' couplings, e.g. in the quark sector, via heavy fermions [Altmannshofer, Gori, Pospelov, Yavin, '14] or Higgs fields. [Crivellin, D'Ambrosio, **JH**, PRL & PRD '15]
- Can get $bs\mu\mu$ operator, explains $R(K)$ and $R(K^*)$!
- LFUV from $L_\mu - L_\tau$, but Z' needs to be heavier than for g-2.
- Predicts deviations in $b \rightarrow s\tau\tau$ and $b \rightarrow sv\nu$, but tough to test.
- Strong constraint from $B-\bar{B}$ oscillations.



$L_\mu - L_\tau$ beyond the Z'

- Can use $L_\mu - L_\tau$ (or other flavor $U(1)$'s) to forbid or enforce couplings.
- Take leptoquark $S_3 \sim (\mathbf{3}, \mathbf{3}, -1/3)$: $\mathcal{L} = y_{ij} \bar{Q}_i S_3 L_j^c + x_{ij} Q_i Q_j S_3$
- Charge $S_3 \sim +1$ under $L_\mu - L_\tau$ to get $x_{ij} = 0$, $y_{ij} = y_{i\mu}$.
→ no more **proton decay**, no **lepton flavor violation**, only coupling to **muons**!
- Perfect for **$R(K)$** & **$R(K^*)$** !
[Hambye, **JH**, PRL '18; Davighi, Kirk, Nardecchia, '20; Greljo, Stangl, Thomsen, '21]
- $L_\mu - L_\tau$ global or gauged, just need to break it in neutrino sector.



Summary

- Standard Model symmetry/prediction:

$$U(1)_{B-L} \times U(1)_{L_\mu - L_\tau} \times U(1)_{L_\mu + L_\tau - 2L_e} .$$

- Easy to gauge any subgroup of this.
- $L_\mu - L_\tau$ special:
 - Only acts on 2nd & 3rd gen particles → weak constraints.
 - Good flavor symmetry for quasi-degenerate **neutrinos**.
 - Only Z' left to explain muon **$g-2$** , soon fully probed.
 - Breaks lepton flavor universality, can explain **$R(K)$** .

Arguably most useful $U(1)$!