

CP-Violating Neutrino Non-Standard Interactions in Long-Baseline-Accelerator Data

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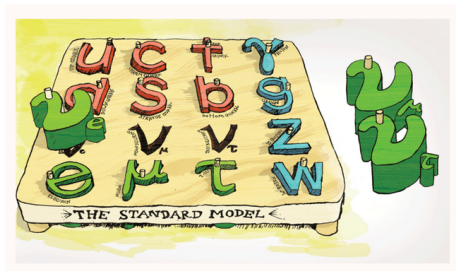
Anomalies 2021

November 2021

Based on 2008.01110, Phys. Rev. Lett. 126 (2021) no.5, 051801
in collaboration with Peter B. Denton and Rebekah Pestes

Neutrino oscillations

- ▶ Neutrino oscillations: **strong** evidence for BSM physics



Neutrino oscillations

- ▶ Neutrino oscillations: strong evidence for BSM physics
- ▶ flavor eigenstates (of weak interaction) and mass eigenstates (of free particle Hamiltonian) not aligned for neutrinos

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

U_{PMNS} : relates flavor and mass states

parametrized by 4 parameters (3 angles, at least 1 phase)

⇒ observation of neutrino oscillation introduced more parameters to the SM

Neutrino oscillations

- ▶ standard parametrization for PMNS matrix as a series of three rotations

$$U_{PMNS} = U_{23}(\theta_{23})U_{13}(\theta_{13}, \delta)U_{12}(\theta_{12})\text{diag}(e^{i\alpha_1/2}, e^{i\alpha_2/2}, 1)$$

- ▶ $\text{diag}(e^{i\alpha_1/2}, e^{i\alpha_2/2}, 1)$ only physical for Majorana neutrinos, oscillation experiments are not sensitive to these phases
→ not going to talk about them further

Neutrino oscillations

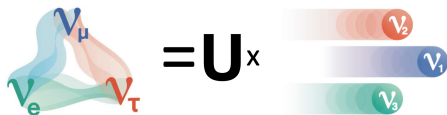
standard parametrization for PMNS matrix as a series of three rotations

$$U_{PMNS} = U_{23}(\theta_{23})U_{13}(\theta_{13}, \delta)U_{12}(\theta_{12})$$

⇒ **want to measure these new parameters!**



Neutrino oscillations

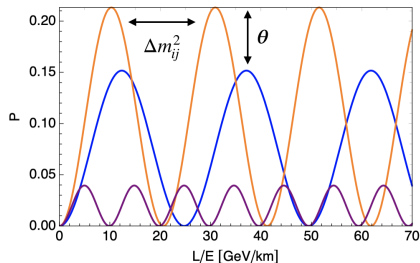


produce neutrino of flavor α with energy E

probability to detect neutrino with flavor β at distance L is

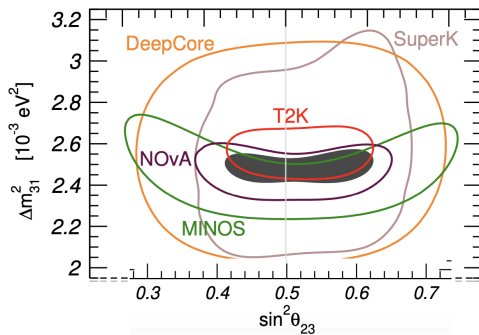
$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2(\Delta m_{ij}^2 L/4E), \quad \Delta m_{ij}^2 = m_i^2 - m_j^2$$

in a 2-flavor approximation



Neutrino oscillations

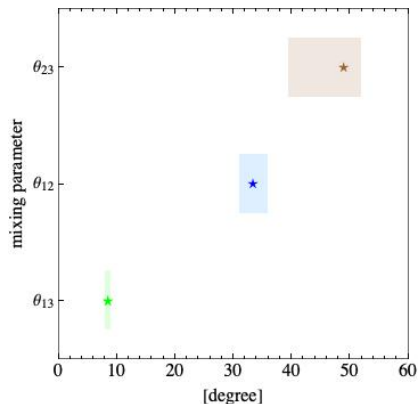
- ▶ many experiments have measured the angles and mass splittings
→ impressive agreement between experiments



[nuFit v5.0]

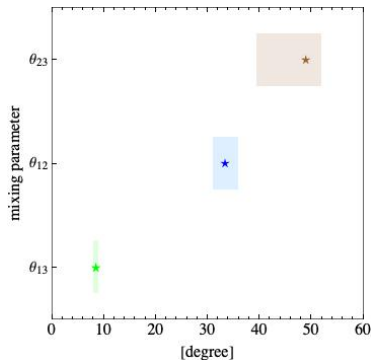
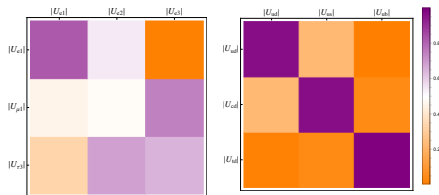
Neutrino oscillations

- ▶ all three angles non-zero
- ▶ mixing angles are large!



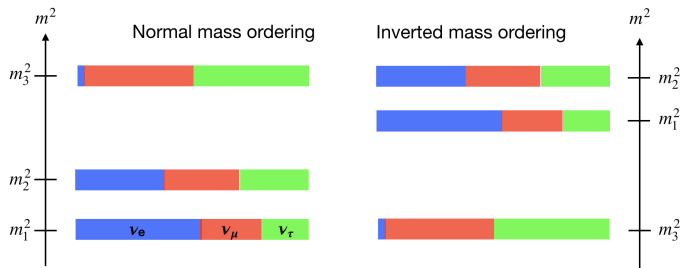
Neutrino oscillations

- ▶ all three angles non-zero
- ▶ mixing angles are large!
surprising if compared to small quark mixing

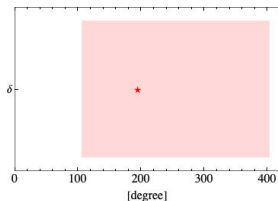


Neutrino oscillations

- ▶ mass splittings: $|\Delta m_{32}^2| = 2.5 \cdot 10^{-3} \text{ eV}^2$, $\Delta m_{21}^2 = 7.4 \cdot 10^{-5} \text{ eV}^2$
- ▶ mass ordering unknown



- ▶ all three mixing angles are non-zero \rightarrow possibility for CPV
- ▶ currently least known parameter is δ which governs CPV in lepton sector
- ▶ want to measure δ !



CPV in the SM

- ▶ weak interaction: CP maximally violated [Cronin, Fitch '64]
- ▶ strong interaction: no observed EDM \rightarrow CP conserved (?) (\rightarrow strong CP problem)



- ▶ CPV in mass matrices quantified via basis invariant

$$J_{CP} = \sin \theta_{13} \cos^2 \theta_{13} \sin \theta_{12} \cos \theta_{12} \sin \theta_{23} \cos \theta_{23} \sin \delta$$

$$J_{CP}^{max} = 1/(6\sqrt{3}) \approx 0.096 \quad [\text{Jarlskog '85}]$$

- ▶ quark mixing matrix: non-zero δ_{CKM} but CPV is small

$$|J_{CKM}|/J_{CP}^{max} = 3 \cdot 10^{-5}$$

- ▶ **is CP violated in the lepton sector?** $|J_{PMNS}|/J_{CP}^{max} < 0.34$

How to measure delta?

- ▶ CPV can only take place in appearance experiments $P(\nu_\alpha \rightarrow \nu_\beta)$
- ▶ need a channel where all three flavors play a role (need interference of two contributions to the oscillation probability given by the two mass splittings)
- ▶ compare neutrino with anti neutrino oscillation probability

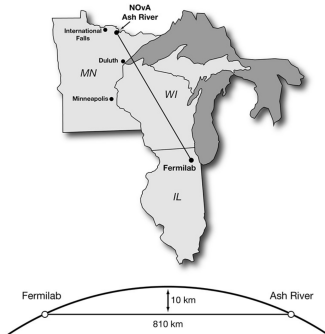
use $P(\nu_\mu \rightarrow \nu_e)$ as oscillation channel!

due to matter effects this channel is also sensitive to the MO

Long baseline experiments: NOvA



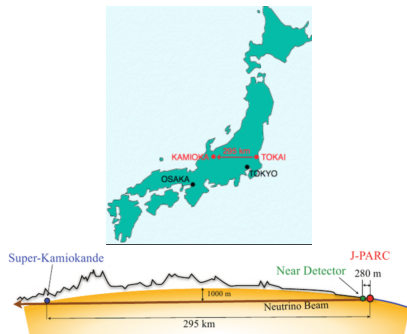
- ▶ neutrinos from NUMI beam at Fermilab
- ▶ $E \sim 1.9$ GeV, $L=810$ km
- ▶ matter density $\rho = 2.84$ g/cc
- ▶ Near (far) detector 0.3 (14) kT liquid scintillator



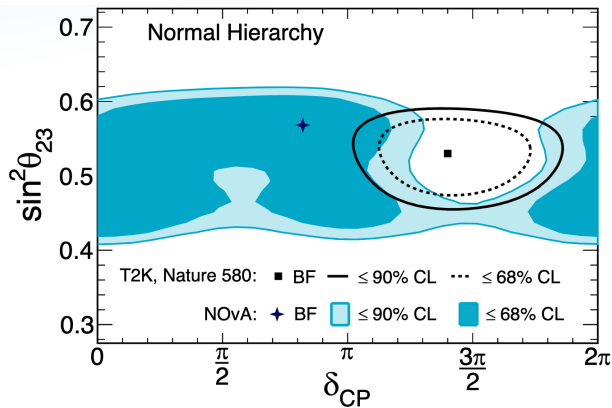
Long baseline experiments: T2K



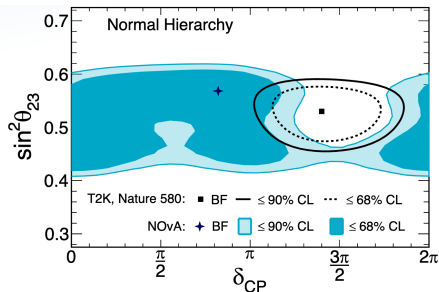
- ▶ neutrinos from JPARC beam
- ▶ $E \sim 0.6$ GeV, $L=295$ km
- ▶ matter density $\rho = 2.3$ g/cc
- ▶ near detector: plastic scintillator, far detector is SuperK



Excitement at Neutrino2020!



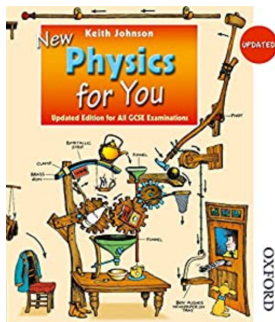
[Himmel '20]



[Himmel '20]

- ▶ both experiments prefer NO
- ▶ no strong preference for NOvA, generally around $\delta \sim \pi$
- ▶ T2K prefers $\delta = 3\pi/2$
 - ⇒ slight disagreement at $\sim 2\sigma$

Can new physics alleviate this slight discrepancy?



difference between NOvA and T2K is the baselines and the matter density

→ neutrinos at NOvA experience stronger matter effects

new physics solution could be related to this difference

introduce new matter interactions for neutrinos

⇒ **neutrino non-standard interactions**

- ▶ framework

$$\mathcal{L}_{NSI} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f} \epsilon_{\alpha\beta}^f (\bar{\nu}_\alpha \gamma^\mu \nu_\beta) (\bar{f} \gamma_\mu f)$$

- ▶ affect oscillations via new matter effect

$$H = \frac{1}{2E} \left[U^\dagger M^2 U + a \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix} \right]$$

matter potential $a \propto G_F \rho E$

focus on vectorial NSI, flavor changing parameters

6 real parameters:

$$|\epsilon_{e\mu}|e^{i\phi_{e\mu}}, |\epsilon_{e\tau}|e^{i\phi_{e\tau}}, |\epsilon_{\mu\tau}|e^{i\phi_{\mu\tau}}$$

consider only one complex parameter at a time

assumption: NSI only affects δ , θ_{23} & Δm_{31}^2 unaffected

$$P(\epsilon = 0, \delta_{\text{meas}}) = P(\epsilon, \delta_{\text{true}}),$$

$$\bar{P}(\epsilon = 0, \delta_{\text{meas}}) = \bar{P}(\epsilon, \delta_{\text{true}}),$$

using approximate expressions for NSI in LBL from [Kikuchi '09]

→ estimates for needed magnitude and phase of NSI parameter

estimate for magnitude of NSI parameter

$$|\epsilon_{e\beta}| \approx \frac{s_{12}c_{12}c_{23}\pi\Delta m_{21}^2}{2s_{23}w_\beta} \left| \frac{\sin\delta_{T2K} - \sin\delta_{NOvA}}{a_{NOvA} - a_{T2K}} \right| \approx \begin{cases} 0.22 & \text{for } \beta = \mu \\ 0.24 & \text{for } \beta = \tau \end{cases},$$

$$\text{with } w_\beta = \sin\theta_{23}(\cos\theta_{23})$$

- ▶ if $\sin\delta_{T2K} = \sin\delta_{NOvA} \rightarrow |\epsilon| = 0$
- ▶ if $a_{T2K} = a_{NOvA} \rightarrow |\epsilon| \rightarrow \infty$

estimate for phase of NSI parameter

with $\delta_{\text{NOVA}} \neq \delta_{\text{T2K}}$

we find that $\sin(\delta_{\text{true}} + \phi_{e\beta}) \approx 0$

with $a_{\text{NOVA}} > a_{\text{T2K}}$ and $\sin \delta_{\text{T2K}} < \sin \delta_{\text{NOVA}}$:

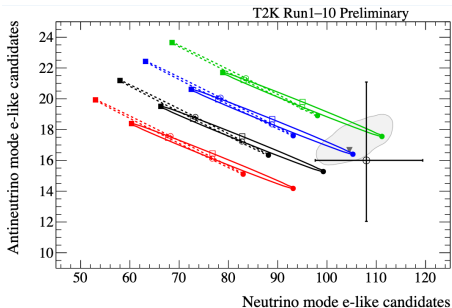
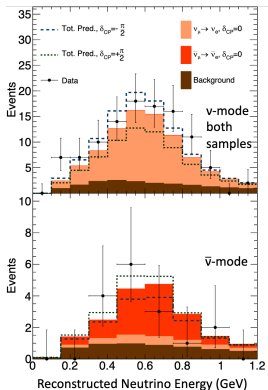
$$\cos(\delta_{\text{true}} + \phi_{e\beta}) \approx -1$$

$$\delta_{\text{true}} \approx \delta_{\text{T2K}} \rightarrow \phi_{e\beta} \approx \frac{3}{2}\pi$$

► Appearance data:

$$n(\nu_e) = xP(\nu_\mu \rightarrow \nu_e) + yP(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) + z,$$

include wrong sign leptons, fit to points on biprobability plot



- ▶ Appearance data:

$$n(\nu_e) = xP(\nu_\mu \rightarrow \nu_e) + yP(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) + z,$$

include wrong sign leptons, fit to points on biprobability plot

- ▶ Disappearance data:

NOvA: use fit results from [Kelly et al '20]

T2K: use provided likelihoods

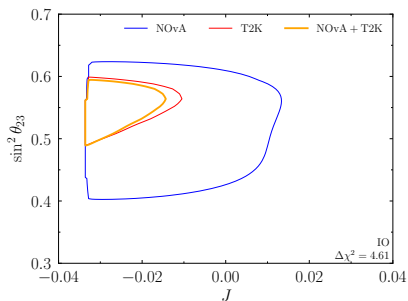
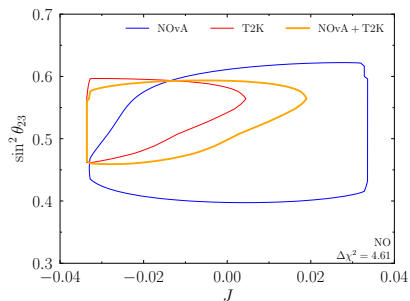
use information from vacuum experiments for remaining parameters

▶ Daya Bay: θ_{13} , $|\Delta m_{32}^2|$ [1809.02261]

▶ KamLAND: θ_{12} , $|\Delta m_{21}^2|$ [1303.4667]

▶ SNO: $\Delta m_{21}^2 > 0$

conduct combined analysis of NOvA and T2K using a log likelihood ratio with Poisson statistics

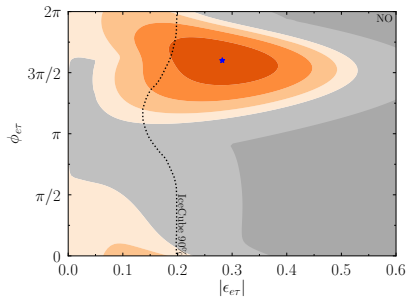
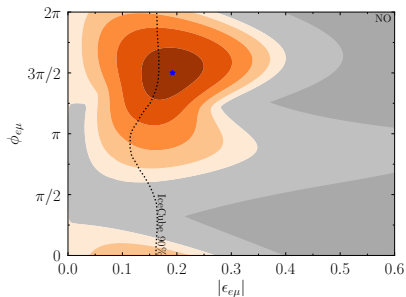


combination is more compatible with IO, IO preferred over NO at $\Delta\chi^2 = 2.3$

discrepancy slightly resolved by swapping the mass ordering

1. NOvA and T2K both prefer NO over IO
2. NOvA+T2K prefers IO over NO
3. SK still prefers NO over IO
4. NOvA+T2K+SK still prefers NO over IO
5. near future reactor experiments provide information in the future

[Kelly et al '20; Esteban et al '20; Denton, JG, Pestes '20]



orange preferred over SM at integer values of $\Delta\chi^2$, dark gray disfavored at $\Delta\chi^2 = 4.61$

[see also Chatterjee, Palazzo '20]

analytical estimates: $|\epsilon_{\alpha\beta}| \approx 0.2$, $\phi_{\alpha\beta} \approx 3\pi/2$, $\delta_{\text{true}} \approx 3\pi/2$

MO	NSI	$ \epsilon_{\alpha\beta} $	$\phi_{\alpha\beta}/\pi$	δ/π	$\Delta\chi^2$
NO	$\epsilon_{e\mu}$	0.19	1.50	1.46	4.44
	$\epsilon_{e\tau}$	0.28	1.60	1.46	3.65
	$\epsilon_{\mu\tau}$	0.35	0.60	1.83	0.90
IO	$\epsilon_{e\mu}$	0.04	1.50	1.52	0.23
	$\epsilon_{e\tau}$	0.15	1.46	1.59	0.69
	$\epsilon_{\mu\tau}$	0.17	0.14	1.51	1.03

$$\Delta\chi^2 = \chi_{\text{SM}}^2 - \chi_{\text{NSI}}^2 \text{ for a fixed MO, } \chi_{\text{NO}}^2 - \chi_{\text{IO}}^2 = 2.3$$

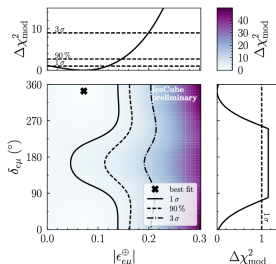
Oscillation constraints on NSI parameters

NSI effects grow with energy, distance and matter density

- ▶ $\epsilon_{\mu\tau}$ best probed with atmospheric neutrinos
- ▶ $\epsilon_{e\mu}, \epsilon_{e\tau}$ best probed with LBL appearance, atmospheric neutrinos

- ▶ IceCube: slightly disfavors LBL best fit points, prefers non-zero $|\epsilon_{e\mu}|$ at 1σ

- ▶ SuperK: only considered real NSI, comparable sensitivity to IceCube



[IceCube '21]

- ▶ CEvNS at COHERENT: applies to $M_{Z'} > 10$ MeV, comparable constraints [Denton, JG '20]

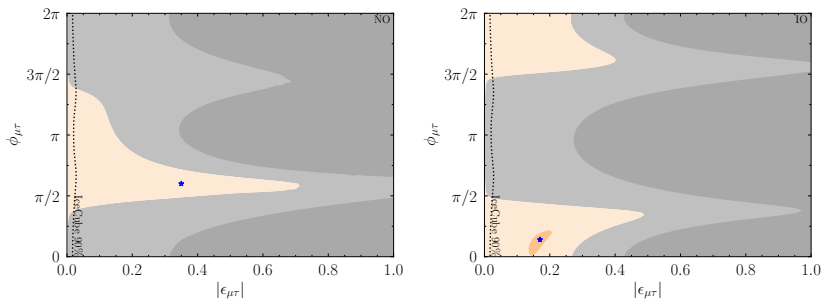
Summary & Conclusions

- ▶ open question of CPV in lepton sector
- ▶ slight tension between NOvA and T2K
- ▶ swap in mass ordering NO \rightarrow IO can resolve this partially
- ▶ tension fully resolved with NSI!
- ▶ predict maximal CP violation in PMNS matrix and for new physics
- ▶ hint for NSI can be further probed with near-future experiments like T2HK and DUNE

Thank you for your attention!



in general no preference for IO with NSI parameters



orange preferred over SM at integer values of $\Delta\chi^2$, dark gray disfavored at $\Delta\chi^2 = 4.61$

$$P(\epsilon = 0, \delta_{\text{meas}}) = P(\epsilon, \delta_{\text{true}}),$$

$$\bar{P}(\epsilon = 0, \delta_{\text{meas}}) = \bar{P}(\epsilon, \delta_{\text{true}}),$$

$$-s_{12}c_{12}c_{23}\frac{\pi}{2}\Delta m_{21}^2 \sin \delta + a_{\text{NOvA}}|\epsilon_{e\beta}| \left[w_{\beta}s_{23} \cos(\delta + \phi_{e\beta}) - v_{\beta}c_{23}\frac{\pi}{2} \sin(\delta + \phi_{e\beta}) \right] \approx -s_{12}c_{12}c_{23}\frac{\pi}{2}\Delta m_{21}^2 \sin \delta_{\text{NOvA}}, \quad (\text{A3})$$

$$s_{12}c_{12}c_{23}\frac{\pi}{2}\Delta m_{21}^2 \sin \delta - a_{\text{NOvA}}|\epsilon_{e\beta}| \left[w_{\beta}s_{23} \cos(\delta + \phi_{e\beta}) + v_{\beta}c_{23}\frac{\pi}{2} \sin(\delta + \phi_{e\beta}) \right] \approx s_{12}c_{12}c_{23}\frac{\pi}{2}\Delta m_{21}^2 \sin \delta_{\text{NOvA}}, \quad (\text{A4})$$

► NSI framework

[Dev et al '19]

$$\mathcal{L}_{NSI} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f} \epsilon_{\alpha\beta}^f (\bar{\nu}_\alpha \gamma^\mu \nu_\beta) (\bar{f} \gamma_\mu f)$$

→ NSI affect neutrino-SM scattering experiments like CEvNS

► NSI effect on weak charge

$$Q_{W\alpha}^2 \propto \left[N(g_n^V + \epsilon_{\alpha\alpha}^V) \right]^2 + \sum_{\beta \neq \alpha} \left[\epsilon_{\alpha\beta}^V (Z + N) \right]^2$$

→ no sensitivity to complex NSI phase

Backup: Scattering constraints on NSI parameters

- ▶ CEvNS detected for the first time in 2016 by COHERENT
- ▶ NSI constraints apply to mediators $M_{Z'} > 10$ MeV
- ▶ constraints derived in [Denton, JG '20] using Feldman-Cousins framework

