Two-Higgs doublet solution to the LSND, MiniBooNE and muon g - 2 anomalies

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Anomalies 2021, IITH

- LSND and MiniBooNE (MB) excesses
- Present constraints on light sterile neutrino
- Decay of heavy sterile neutrino as a solution
- Constraints on our model
- Conclusions

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LSND and MB excess

- LSND and MB are two short-baseline neutrino experiments.
- Schematic representation of MB :



- \bullet LSND : proton energy 800 MeV and base-line \sim 30 meters.
- Detectors can't distinguish the signals from e^-, e^+ , and γ .

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MB and LSND fluxes:



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• MB looked for an *electron like* signal in the final states.

 ν_{μ} beam and $\nu_{\mu} + A \rightarrow e^{-} + X$

• LSND signal looks like an inverse β -decay.

 $\bar{\nu}_{\mu}$ beam and $\bar{\nu}_{e} + p \rightarrow e^{+} + n$

and n is captured by the free hydrogen in the detector. This produces a unique signature of 2.2 MeV gamma in the detector.

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MB and LSND events



• An excess of *electron (positron) like* events over the expected background was observed.

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- Solution to the LSND and MB anomalies via light sterile neutrino does not fit very well in the global picture.
- Recent results of MicroBooNE also disfavor the light sterile neutrino hypothesis. (arXiv : 2110.14065)

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Some general constraints

MB beam dump run :



- If events scale as POT, as in DM production and scattering, then 35.5 excess events expected.
- However, only 6 events were seen, when expected background was 8.8.

Conclusion : Excess disappears when neutrino flux is suppressed, and is thus linked to neutrinos.

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The proposed model:

We consider two Higgs doublet model (2HDM) with a dark singlet scalar $\phi_{h'}$. In addition, three right-handed neutrinos help to generate neutrino masses via the seesaw mechanism and participate in the interaction which gives the signal in MB and LSND.

$$\begin{split} V &= |\phi_h|^2 \left(\frac{\lambda_1}{2} |\phi_h|^2 + \lambda_3 |\phi_H|^2 + \mu_1\right) \qquad \text{where} \\ &+ |\phi_H|^2 \left(\frac{\lambda_2}{2} |\phi_H|^2 + \mu_2\right) + \lambda_4 (\phi_h^{\dagger} \phi_H) (\phi_H^{\dagger} \phi_h) \qquad \qquad \phi_h = \left(\frac{H_1^+}{\frac{\nu + H_1^0 + iG^0}{\sqrt{2}}}\right), \\ &+ \phi_{h'}^2 \left(\lambda_2' \phi_{h'}^2 + \lambda_3' |\phi_h|^2 + \lambda_4' |\phi_H|^2 + m' \phi_{h'} + \mu'\right) \\ &+ \left[\phi_h^{\dagger} \phi_H \left(\frac{\lambda_5}{2} \phi_h^{\dagger} \phi_H + \lambda_6 |\phi_h|^2 + \lambda_7 |\phi_H|^2 + \lambda_5' \phi_{h'}^2 - \mu_{12}\right) \qquad \phi_H = \left(\frac{H_2^+}{\frac{H_2^0 + iA^0}{\sqrt{2}}}\right) \\ &+ \phi_{h'} (m_1 |\phi_h|^2 + m_2 |\phi_H|^2 + m_{12} \phi_h^{\dagger} \phi_H) + h.c.\right] \qquad \phi_{h'} = H_3^0 / \sqrt{2} \end{split}$$

In the Higgs basis the relevant Lagrangian ${\cal L}$ can be written as

$$\begin{split} \mathcal{L} &= \sqrt{2} \Big[(X^{u}_{ij} \tilde{\phi}_{h} + \bar{X}^{u}_{ij} \tilde{\phi}_{H}) \bar{Q}^{i}_{L} u^{j}_{R} + (X^{d}_{ij} \phi_{h} + \bar{X}^{d}_{ij} \phi_{H}) \bar{Q}^{i}_{L} d^{j}_{R} \\ &+ (X^{e}_{ij} \phi_{h} + \bar{X}^{e}_{ij} \phi_{H}) \bar{L}^{i}_{L} e^{j}_{R} + (X^{\nu}_{ij} \tilde{\phi}_{h} + \bar{X}^{\nu}_{ij} \tilde{\phi}_{H}) \bar{L}^{i}_{L} \nu_{R_{j}} \\ &+ \frac{1}{\sqrt{8}} m_{ij} \bar{\nu}^{c}_{R_{i}} \nu_{R_{j}} + \lambda^{N}_{ij} \bar{\nu}^{c}_{R_{i}} \phi_{h'} \nu_{R_{j}} + h.c. \Big] \end{split}$$

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Neutrino interactions:

$$\mathcal{L}_{\nu}^{\text{int}} \simeq y_{\nu_{ij}}^{\phi} \bar{\nu}_i N_j \phi + (\lambda_{N_{ij}}^{h'} h' + \lambda_{N_{ij}}^H H) \bar{N}_i N_j + h.c.$$
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m_{N_1}	m_{N_2}	m_{N_3}	$y_u^{h'(H)} \times 10^6$	$y_{e(\mu)}^{h'} \times 10^4$	$y_{e(\mu)}^{H} \times 10^{4}$
$85 \mathrm{MeV}$	$130\mathrm{MeV}$	$10{ m GeV}$	0.8(8)	0.23(1.6)	2.29(15.9)
$m_{h'}$	m_H	$\sin \delta$	$y_d^{h'(H)} \times 10^6$	$y_{\nu_{i2}}^{h'(H)} \times 10^3$	$\lambda_{N_{12}}^{h'(H)} \times 10^3$
$17 \mathrm{MeV}$	$750\mathrm{MeV}$	0.1	0.8(8)	1.25(12.4)	74.6(-7.5)

TABLE I: Benchmark point used for event generation in LSND, MB and for calculating the muon g - 2.

Results:



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Discussion of some relevant points:

• All LSND events in our scenario stem from the high energy part of the DIF flux, which is kinematically capable of producing N_2 ($m_{N_2} = 130$ MeV).

$$\nu_{\mu} CH_2 \rightarrow nN_2 X \rightarrow nN_1 h' X \rightarrow N1 \gamma e^+ e^- X$$

- The masses of N_2 and N_1 are important factors in obtaining both the correct number and the correct distributions of events in the detectors.
- KARMEN, a smaller experiment similar to LSND did not observe any excess. N₂ will not be produced as the high energy DIF flux of KARMEN is very small.



CHARM II and MINER $\nu {\rm A}$ constrain the proposed model by the $\nu_{\mu}-e$ scattering data.

 $\nu_{\mu} A \rightarrow N_2 A$ coherent cross section becomes relevant. N_2 will decay and produce an electron-like signal in these detectors, and potentially conflict with measured $\nu - e$ data.

 $\nu_{\mu}\,A \rightarrow \textit{N}_{2}\,A$ coherent cross section remains below the 1% of $\nu-e$ cross section in our model.

IceCube provides constraints on our model via the searches of double bang events. The decay time of N_2 (leading to e^+e^- pair) is short enough, to escape the detection at this detector. The distances traveled even at very high energies are much less smaller than the resolution necessary to signal a double bang events, ~ 1 m in DeepCore, and \sim a few hundred meters in IceCube.

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- Evidence for anomalous signals at low-energy, and short-baseline neutrino experiments in particular, increased over time.
- Our proposed model could provide a common, non-oscillation, new physics explanation for both LSND and MB. It also explains the anomalous magnetic moment of muon.
- Our model does not conflict with the recent results of MicroBooNE.

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Thank you for your attention

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