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# Fermi Surface Studies Of Co-based Heusler Alloys: Ab-initio Study

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**Abstract.** The electronic, Fermi surface (FS) and magnetic properties of ferromagnetic Heusler alloys  $\text{Co}_2\text{XY}$  ( $X = \text{Cr}, \text{Mn}, \text{Fe}; Y = \text{Al}, \text{Ga}$ ) have been investigated by means of first principles calculation. Out of these compounds,  $\text{Co}_2\text{CrAl}$  is found to be perfectly half-metallic (HM) at ambient. Under pressure HM to nearly HM (NHM) transition is observed around 75 GPa for  $\text{Co}_2\text{CrAl}$  and NHM to HM transition is observed around 40 GPa and 18 GPa for  $\text{Co}_2\text{CrGa}$  and  $\text{Co}_2\text{MnAl}$ , respectively, while no transition is observed for other compounds under study and is also analyzed from the FS studies. The states at the Fermi level in the majority spin are strongly hybridized Co-d and X-d like states. The majority band FS topology change is observed under pressure for the compounds where we observe a transition, while the minority band FS remain unaltered under pressure for all compounds except in  $\text{Co}_2\text{FeGa}$ , where we observed an electron sheet at X point instead of hole pocket at  $\Gamma$  point.

**Keywords:** Magnetic properties and materials

**PACS:** 75

## INTRODUCTION

The Co-based Heusler alloys have a great attention due to their high Curie temperature [1]. Theoretically it was predicted that most of Co-based Heusler alloys behave like HM, where they are expected to provide 100% spin-polarization at Fermi level ( $E_F$ ) [2]. These half-metallic ferromagnetic materials can be used as a perfect spin filters or to enhance the performance of spin-dependent devices. In spite of these, there are many discrepancies in the electronic properties of these compounds. Ishida et. al [3] have shown that  $\text{Co}_2\text{CrAl}$  has a complete spin polarization at the  $E_F$  which contradicts the prediction of Galanakis [4].  $\text{Co}_2\text{XY}$  type Heusler alloys are good candidates for spintronics materials, where X atoms belong to transition metal and Y is a main group element.

## METHOD OF CALCULATION

We have carried out density functional calculations using the full potential linear augmented plane wave method. For the exchange-correlation functional, the generalized gradient approximation (GGA) was used. The self-consistent calculations were considered to be converged when the total energy of the system was stable within  $10^{-6}$  Ry. A  $(44 \times 44 \times 44)$  mesh was used

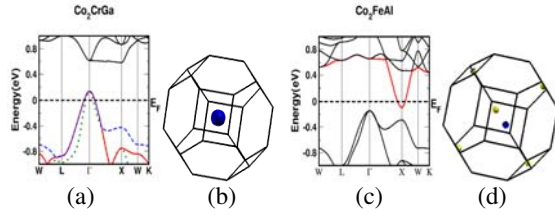
for the Fermi surface calculations to ensure accurate determination of the  $E_F$ .

## Results And Discussion (Band Structure, Fermi Surface And Density Of States)

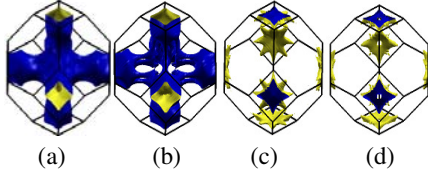
We have calculated the density of states ( $N(E_F)$ ) for minority ( $\downarrow$ ) and majority ( $\uparrow$ ) spin and the spin polarization (SP) ratio at  $E_F$  and are reported in TABLE 1, which confirm  $\text{Co}_2\text{CrAl}$  to be a perfect HM with 100% SP at  $E_F$ , while  $\text{Co}_2\text{CrGa}$ ,  $\text{Co}_2\text{MnAl(Ga)}$ ,  $\text{Co}_2\text{FeAl(Ga)}$  are NHM with SP less than 100%. The observed NHM nature of these compounds is mainly due to the valence band maximum ( $\text{VB}_{\text{max}}$ ) crossing the  $E_F$  and forming three hole pockets at the  $\Gamma$  point for corresponding three bands (FIGURE 1(a, b)) which are Co-d ( $t_{2g}$ ) like states except in  $\text{Co}_2\text{FeAl}$ , where we found the conduction band minimum ( $\text{CB}_{\text{min}}$ ) to cross the  $E_F$  at X point, thereby forming an electron pocket (FIGURE 1(c, d)) and is mainly Fe-d( $e_g$ ) like states. The presence of the hole or electron pocket for minority spin band with finite value of  $N(E_F)$  as reported in TABLE 1, further confirm the NHM behavior of these compounds, except  $\text{Co}_2\text{CrAl}$ , which is completely HM with zero  $N(E_F)$  for the minority band. The majority spin band of all compounds are strongly metallic in nature and are hybridized Co-d

and X-d orbitals with small contribution from Al(Ga)-p like states and is again reflected in the FS topology. The FS topology remains unchanged by replacing Al with Ga, which indicates the minimum contribution of Al(Ga) at  $E_F$ , but the FS topology do change when we vary the X(Cr, Mn, Fe) atom (FIGURE 2(a, c)).

**FIGURE 1.** (a, b) Minority spin band structure and FS of  $\text{Co}_2\text{CrGa}$ . The topology of other two surfaces is also same as 1<sup>st</sup> surface. In addition the topology for  $\text{Co}_2\text{MnAl}$ (Ga) and  $\text{Co}_2\text{FeGa}$  is also same as  $\text{Co}_2\text{CrGa}$ ; (c, d) Minority spin band structure and FS of  $\text{Co}_2\text{FeAl}$ .



**FIGURE 2.** Majority spin FS, (a, c) for  $\text{Co}_2\text{CrAl}$  and  $\text{Co}_2\text{MnAl}$  at ambient; (b, d) under pressure. The topology of Al and Ga varying compounds are same.

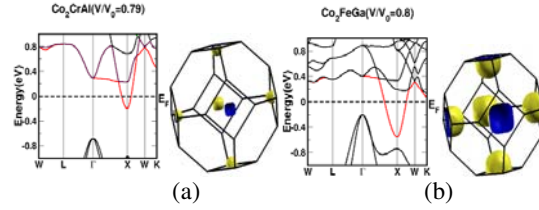


Under pressure at nearly 75GPa, the HM behavior of  $\text{Co}_2\text{CrAl}$  is destroyed due to the downward shift of the band at X point, resulting in an electron pocket at X point in minority spin as shown in FIGURE 3(a). Similarly around 18GPa and 48GPa,  $\text{Co}_2\text{CrGa}$  and  $\text{Co}_2\text{MnAl}$  behave like a HM due to the downward shift of the band at  $\Gamma$  point, resulting in the vanishing of hole pocket for the minority spin band. The NHM nature of the other investigated compounds remains unaltered and the FS topology is also found to be the same under pressure. In addition to this, we observed an electron pocket at X point instead of hole pocket at  $\Gamma$  for  $\text{Co}_2\text{FeGa}$  due to the downward shifts of the band which is mainly Fe-d( $e_g$ ) like states and is evident from FIGURE 3(b). The main interesting feature of these investigated compounds is the observed FS topology change under pressure.

**TABLE 1.** Calculated  $N(E_F)$  for majority ( $\uparrow$ ) and minority ( $\downarrow$ ) spin in (states/eV/spin) at the  $E_F$ , gap between  $CB_{\min}$  and  $VB_{\max}$ , spin polarization (SP) in % at the  $E_F$ , total magnetic moment in  $\mu_B$ . All the calculated values are at zero pressure.

Compounds	$N(E_F\uparrow)$	$N(E_F\downarrow)$	SP	Gap in eV	$\mu_B(\text{Total})$
$\text{Co}_2\text{CrAl}$	4.12	0	100	0.73	3.00
$\text{Co}_2\text{CrGa}$	5.42	0.13	95.32	0.44	3.03
$\text{Co}_2\text{MnAl}$	1.02	0.15	74.64	0.67	4.04
$\text{Co}_2\text{MnGa}$	1.62	0.36	63.64	0.38	4.11
$\text{Co}_2\text{FeAl}$	0.87	0.08	83.94	0.05	4.99
$\text{Co}_2\text{FeGa}$	0.88	0.14	72.44	0.02	5.02

**FIGURE 3.** Variation of minority Band and FS under pressure, (a)  $\text{Co}_2\text{CrAl}$  (b)  $\text{Co}_2\text{FeGa}$



We observe the majority band FS topology change for the compounds where transition is seen (HM to NHM in  $\text{Co}_2\text{CrAl}$  and NHM to HM in  $\text{Co}_2\text{CrGa}$  and  $\text{Co}_2\text{MnAl}$ ) as shown in FIGURE 2. Though the topology of  $\text{Co}_2\text{MnAl}$  and  $\text{Co}_2\text{MnGa}$  are the same at ambient, the topology remains unaltered in  $\text{Co}_2\text{MnGa}$ , where we could not observe the magnetic transition till high pressure unlike  $\text{Co}_2\text{MnAl}$ .

## CONCLUSION

Among all the compounds  $\text{Co}_2\text{CrAl}$  is found to be HM and turns NHM under pressure. The non-vanishing  $N(E_F)$  and the minority band FS confirm the NHM behavior of all these investigated compounds. The FS topology change for majority band is observed for the compounds where the transition from HM to NHM or NHM to HM state is observed.

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